# Affine Weyl groups and non-abelian discrete systems

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# **Motivation & Outline**

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- Painlevé equations being one of the most important objects in mathematics and mathematical physics have various types of non-abelian analogs: quantum [Nagoya, 2004], matrix differential [Kawakami, 2015], matrix difference [Cassatella-Contra et al., 2014], non-abelian differential [Bobrova and Sokolov, 2023b].
- Some of them are connected with integrable non-abelian PDEs [Olver and Sokolov, 1998] and P∆Es [Adler, 2020], Riemann-Hilbert problem [Cafasso and Manuel, 2014], orthogonal polynomials [Cafasso et al., 2018], Calogero systems [Bertola et al., 2018], and etc.
- In the commutative case, discrete Painlevé equations have been studied in a series of papers by B. Gramaticos and A. Ramani since 1990s, but without understanding the whole picture.
- The latter was clarified by H. Sakai in his famous paper [Sakai, 2001].
- We would like to derive the same picture in the non-commutative case.
- But we first present an algebraic tool in order to obtain good examples for the further study.
- It uses the affine Weyl groups and might be regarded as a non-abelian analog of that suggested in [Noumi and Yamada, 1998].

#### Outline

- 1. A brief introduction to the Painlevé equations and their non-abelian analogs.
- 2. Affine Weyl groups and discrete dynamics: commutative and non-commutative cases.
- 3. Non-abelian dressing chain and related discrete systems.
- 4. Non-abelian difference discrete Painlevé equations.
- 5. Further questions.

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# Painlevé equations (1)



Number systems and function classes [Joshi, 2020]

- Problem: define new functions by an ODE of the *m*<sup>th</sup> order with properties that generalize those of elliptic functions. [L. Fuchs], [H. Poincaré]
- ▶ Painlevé property: the general solution of an ODE has no critical movable points.
- ▶ m = 2: six classes defining the Painlevé transcendents. [Painlevé, 1902], [Gambier, 1910]
- ▶ P<sub>1</sub> transcendent:  $y''(z) = 6y(z)^2 + z$ . (a non-autonomous analog for the *p*-function!)

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# Painlevé equations (2)



Sakai's classification scheme: the surface type [Sakai, 2001]

(picture source: [Dzhamay and Takenawa, 2018])

- Problem: classify discrete Painlevé equations by using a compact rational surface X with a unique canonical divisor D of canonical type.
- Main idea: to each X, there corresponds a root subsystem of  $E_8^{(1)}$  inside a Picard lattice of X.
- Classification: 22 discrete systems of elliptic, multiplicative, or additive types.
- ► d-P<sub>1</sub>(*E*<sub>7</sub>):  $T(q, p, t; \alpha_0, \alpha_1) = (-q \alpha_1 p^{-1}, -p + 2\bar{q}^2 + t, t; \alpha_0 1, \alpha_1 + 1) \xrightarrow{\text{cont.lim.}} P_1$  $T \in \widetilde{W}(A_1^{(1)})$  and the group is formed by Bäcklund transformations for the P<sub>2</sub>.

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## Why non-abelian?

#### The Ablowitz-Ramani-Segur conjecture [Ablowitz et al., 1980]

A nonlinear PDE is solvable by the inverse scattering method [Zakharov and Shabat, 1974] only if every nonlinear ODE obtained by an exact reduction has the Painlevé property.

► Many important integrable PDEs and P∆Es can be solved in terms of the Painlevé equations. Examples.

- $\blacktriangleright \ \mathsf{KdV} \to \mathsf{P_1}, \, \mathsf{P_2}$
- $\blacktriangleright \ sin-Gordon \to \mathsf{P}_3$
- $\blacktriangleright \ \mathsf{NLS} \to \mathsf{P}_4$
- $\blacktriangleright \ \mathsf{VL} \to \mathrm{dP}_1$
- ► These integrable systems have been intensively studied in the non-commutative setting.
- ▶ Their analogs have similar integrability property generalized to the non-commutative case.
- So, it is natural to investigate their "solutions".

#### A matrix KdV equation [Wadati and Kamijo, 1974]

 $w_t + 6 w_{x} + 6 w_{x} w + w_{xxx} = 0, \qquad w = w(x, t) \in \operatorname{Mat}_n(\mathbb{C}), \qquad x, t \in \mathbb{C}. \quad \operatorname{KdV}^0$ 

- The inverse scattering method [Wadati and Kamijo, 1974].
- A hierarchy of commuting symmetries [Olver and Sokolov, 1998], [Olver and Wang, 2000].
- The ZCR  $\partial_t U \partial_x V = [V, U]$ , where  $U = U(\mu, x, t)$  and  $V = V(\mu, x, t)$ .
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From affine Weyl groups to discrete dynamics

### Affine Weyl groups

- Let us fix a generalized Cartan matrix  $C = (c_{ij})$ , where  $i, j \in I := \{0, 1, \dots, n\}$ .
- Sets  $\Delta = \{\alpha_0, \dots, \alpha_n\}$ ,  $\Delta^{\vee} = \{\alpha_0^{\vee}, \dots, \alpha_n^{\vee}\}$  correspond to simple roots and simple co-roots.
- ▶ Denote by Q = Q(C) and  $Q^{\vee} = Q^{\vee}(C)$  the root and co-root lattices. The pairing  $\langle \cdot, \cdot \rangle : Q \times Q^{\vee} \to \mathbb{Z}$  is defined by  $\langle \alpha_i, \alpha_j^{\vee} \rangle = c_{ij}$  and  $\alpha_i^{\vee} = 2\alpha_i / (\alpha_i, \alpha_i)$ .
- ▶ Denote by W = W(C) the Weyl group (or the Coxeter group) defined by generators  $s_i$ ,  $i \in I$ :

$$W(C) = \langle s_0, s_1, \dots, s_n \mid s_i^2 = 1, (s_i s_j)^{m_{ij}} = 1 \rangle,$$
(1)

where the exponents are determined by the value of the product  $c_{ij}c_{ji}$  as below

$$\begin{array}{c|c} c_{ij}c_{ji} & 0 & 1 & 2 & 3 & \geq 4 \\ \hline m_{ij} & 2 & 3 & 4 & 6 & \infty \end{array}$$

These generators act naturally on Q by reflections

$$\mathbf{s}_i(\alpha_j) = \alpha_j - \langle \alpha_i, \alpha_j^{\vee} \rangle \, \alpha_i = \alpha_j - \mathbf{c}_{ij} \, \alpha_i. \tag{2}$$

- Each s<sub>i</sub>-action on Q induces an automorphism of the C(α) of rational functions in α<sub>i</sub>. Hence, C(α) is a left W-module.
- Recall that one of the important properties of the affine Weyl groups is that they have translations, also known as Kac translations. Let  $W_0$  be a finite Weyl group,  $\delta = \sum_{i \in I} k_i \alpha_i$  be the null root and  $V_0 = \{\mu \in V \mid \langle \mu, \delta^{\vee} \rangle = 0\}$ . For an element  $\mu \in V_0$  such that  $\langle \mu, \mu^{\vee} \rangle \neq 0$  we define a translation element  $t_{\mu} \in W$  by the formula

$$t_{\mu} = s_{\delta - \mu} \, s_{\mu} \tag{3}$$

and suppose that  $w t_{\mu} = t_{w(\mu)} w$  for any  $w \in W$ .

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### Extended birational representations [Noumi and Yamada, 1998]

► The Kac translation acts on simple affine roots as follows

$$t_{\mu}(\alpha) = \alpha - \langle \mu, \alpha \rangle \,\delta = \alpha - \mu_{\alpha} \delta. \tag{4}$$

- It is known that the affine Weyl group is decomposed into a semi-direct product of translations in the lattice part M and the finite Weyl group W<sub>0</sub> acting on M, i.e. W = M ⋊ W<sub>0</sub>. The lattice part M acts on C(α) as a shift operator, thanks to (4).
- Let us consider the set of functions  $f_i$ ,  $i \in I$ , which we will often call variables.
- We propose an extension of the representation of W on  $\mathbb{C}(\alpha)$  to the field  $\mathbb{C}(\alpha, f)$  of rational functions in  $\alpha_i$  and  $f_i$ ,  $i \in I$ . One needs to specify the action of  $s_i$  on  $f_j$  in such a way that the automorphisms  $s_i$  on  $\mathbb{C}(\alpha, f)$  preserve the Weyl group structure.

**Remark 1.** Such classes of representations arise naturally from Bäcklund transformations of the differential Painlevé equations.

**Remark 2.** Sometimes it is necessary to work with an extended Weyl group W, which is a semi-direct product of W and the group  $\Omega$  of automorphisms of the Dynkin diagram  $\Gamma(C)$ , i.e.

$$\widetilde{W} = W \rtimes \Omega.$$

An automorphism of  $\Gamma(C)$  is a bijection  $\pi$  on I s.t.  $c_{\pi(i)\pi(j)} = c_{ij}$  and, therefore,  $\pi s_i = s_{\pi(i)} \pi$ . Note that the representations of W lifts to a representation of  $\widetilde{W}$ .

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## Discrete dynamics [Noumi and Yamada, 1998]

- Suppose that we extended the action of W from  $\mathbb{C}(\alpha)$  to  $\mathbb{C}(\alpha, f)$ . Here we consider an arbitrary extension  $\mathbb{C}(\alpha, f)$  as a left W-module, assuming that each element of W acts on the function field as an automorphism.
- ▶ For each  $\mu \in M$  we define a set of rational functions  $F_{\mu,i}(\alpha, f) \in \mathbb{C}(\alpha, f)$  by

$$t_{\mu}(f_i) = F_{\mu,i}(\alpha, f). \tag{5}$$

- This set can be considered as a discrete dynamical system.
- Here  $\alpha_i$  and  $f_i$  are discrete time variables and depended variables respectively.
- Based on the action of t<sub>µ</sub> on the discrete time variables, discrete dynamics can be classified into additive (*d*-equations), multiplicative (*q*-equations), or elliptic (*ell*-equations) types.

**Remark.** A similar description of the discrete dynamics can be given for  $\widetilde{W}$  as well.

Now it is clear how to generalise this construction to the non-abelian case. We just need to consider "non-commutative" f<sub>i</sub> and repeat the same arguments as above.

# Non-abelian setting

- $\blacktriangleright$  Consider an associative unital division ring  $\mathcal R$  over the field  $\mathbb C$  equipped with a derivation.
- We assume that all greek letters belong to the field  $\mathbb{C}$ , while the elements  $f_i$  are from  $\mathcal{R}$ . We will often call  $f_i$  as *functions*.
- ▶ The derivation  $d_t : \mathcal{R} \to \mathcal{R}$  of the ring  $\mathcal{R}$  is a  $\mathbb{C}$ -linear map satisfying the Leibniz rule. We also assume that there is a central element t such that  $d_t(t) = 1$  and for any  $\alpha \in \mathbb{C}$  we have  $d_t(\alpha) = 0$ . Here and below we identify the unit of the field with the unit of the ring.
- For the brevity we denote  $d_t(f_i) = \dot{f}_i$ ,  $d_t^2(f_i) = \ddot{f}_i$ , and so on.
- Note that on R we have an involution called the transposition τ, which acts trivially on the generators of R and for any elements F, G ∈ R we have τ(F G) = τ(G) τ(F). This involution can be naturally extended to the matrices over R.

**Remark.** We would rather not specify the generators of the ring  $\mathcal{R}$  in order to avoid overloaded description of a pretty simple thing. Instead, we encourage the reader to think of the ring  $\mathcal{R}$  as a generalization of rational functions over the field  $\mathbb{C}$  to a non-abelian case.

## Non-abelian discrete dynamics [Bobrova, 2024]

- Consider the set of elements  $f_i \in \mathbb{R}$ ,  $i \in I$ , which we will often call *functions* or *variables*.
- ▶ Repeating all the arguments above, we again suppose that the action of *W* is extended from  $\mathbb{C}(\alpha)$  to  $\mathbb{C}(\alpha, f)$  and for each  $\mu \in M$  we define a set of elements  $F_{\mu,i}(\alpha, f) \in \mathbb{C}(\alpha, f)$  by

$$t_{\mu}(f_i) = F_{\mu,i}(\alpha, f).$$
(6)

This set can be considered as a non-abelian discrete dynamical system.

**Remark.** Up to the author's knowledge, examples of non-abelian discrete systems of *ell*-type have not appeared yet, while there exist examples of q- and d-types systems. Systems of q-type might be found in [Bobrova et al., 2023], where the non-commutative analogs of the q-P<sub>1</sub> and q-P<sub>2</sub> hierarchies are presented. Examples of non-abelian discrete d-systems can be found, for instance, in [Cassatella-Contra et al., 2014] and [Adler, 2020].

- A suitable birational representation of W leads to a discrete dynamical system.
- Note that classes of such representations arise naturally from Bäcklund transformations of ordinary differential equations, in particular, of the Painlevé equations.

Dynamics related to the dressing chain

### Non-abelian dressing chains

**Remark.** The commutative dressing chain was introduced in [Veselov and Shabat, 1993]. It is related to the Painlevé equations [Adler, 1994] and arises from a generalisation of the symmetries for the P<sub>4</sub> and P<sub>5</sub> systems [Noumi and Yamada, 2000]. Quantum dressing chain might be found in [Nagoya, 2004]. Here we do not assume any relations on the elements  $f_i$ .

▶ Let 
$$j \in \mathbb{Z}/(n+1)\mathbb{Z}$$
. Consider the systems for  $n = 2l$  and  $n = 2l + 1$ ,  $l \in \mathbb{Z}_{\geq 0}$ , respectively

$$\begin{split} \dot{f_j} &= \sum_{1 \le r \le l} f_j \, f_{j+2r-1} - \sum_{1 \le r \le l} f_{j+2r} \, f_j + \alpha_j; \\ \frac{1}{2} t \, \dot{f_j} &= \sum_{1 \le r \le s \le l} f_j \, f_{j+2r-1} \, f_{j+2s} - \sum_{1 \le r \le s \le l} f_{j+2r} \, f_{j+2s+1} \, f_j \\ &+ \left( \frac{1}{2} - \sum_{1 \le r \le l} \alpha_{j+2r} \right) f_j + \alpha_j \sum_{1 \le r \le l} f_{j+2r}. \end{split}$$

• We will cal them  $A_n^{(1)}$ ,  $n \ge 2$  type systems or dressing chains in the Noumi-Yamada variables.

These systems admit Lax pairs.

### Lax pairs

• Let  $\Psi = \Psi(\lambda, t) \in Mat_{n+1}(\mathbb{R})$ ,  $\lambda \in \mathbb{Z}(\mathbb{R})$  satisfy the linear system

$$\begin{cases} \partial_{\lambda}\Psi(\lambda, t) = \mathcal{A}(\lambda, t)\Psi(\lambda, t), \\ \partial_{t}\Psi(\lambda, t) = \mathcal{B}(\lambda, t)\Psi(\lambda, t), \end{cases}$$
(7)

where matrices  $\mathcal{A} = \mathcal{A}(\lambda, t)$  and  $\mathcal{B} = \mathcal{B}(\lambda, t)$  belong to  $Mat_{n+1}(\mathcal{R})$  and are of the form

$$\mathcal{A}(\lambda) = A_0 + A_{-1} \lambda^{-1}, \qquad \qquad \mathcal{B}(\lambda) = B_1 \lambda + B_0. \tag{8}$$

• Consider the matrices expressed in terms of the standard unit matrices  $E_{r,s} \in Mat_{n+1}(\mathbb{C})$  as

$$A_{0} = E_{1,n} + f_{0} E_{1,n+1} + E_{2,n+1}, \quad A_{-1} = \sum_{1 \le r \le n+1} \beta_{r} E_{r,r} + \sum_{1 \le r \le n} f_{r} E_{r+1,r} + \sum_{1 \le r \le n-1} E_{r+2,r},$$

$$B_1 = E_{1,n+1}, \quad B_0 = \sum_{1 \le r \le n+1} g_r E_{r,r} + \sum_{1 \le r \le n} E_{r+1,r}.$$

• Let  $\alpha_0 = 1 + \beta_{n+1} - \beta_1$ ,  $\alpha_j = \beta_j - \beta_{j+1}$ ,  $j \in \mathbb{Z}/(n+1)\mathbb{Z} \setminus \{0\}$ .

**Theorem.** [Bobrova, 2024] There exists a set of the *g*-functions such that the compatibility condition of system (7) is equivalent to either the  $A_{21}^{(1)}$  or  $A_{21+1}^{(1)}$  system.

For the  $A_{2i}^{(1)}$ , we have  $g_j = -\sum_{1 \le r \le l} f_{j+2r}$ , where indexes belong to  $\mathbb{Z}/(n+1)\mathbb{Z}$ .

### Lax pairs

• Let  $\Psi = \Psi(\lambda, t) \in Mat_{n+1}(\mathbb{R})$ ,  $\lambda \in \mathbb{Z}(\mathbb{R})$  satisfy the linear system

$$\begin{cases} \partial_{\lambda}\Psi(\lambda, t) = \mathcal{A}(\lambda, t)\Psi(\lambda, t), \\ \partial_{t}\Psi(\lambda, t) = \mathcal{B}(\lambda, t)\Psi(\lambda, t), \end{cases}$$
(7)

where matrices  $\mathcal{A} = \mathcal{A}(\lambda, t)$  and  $\mathcal{B} = \mathcal{B}(\lambda, t)$  belong to  $Mat_{n+1}(\mathcal{R})$  and are of the form

$$\mathcal{A}(\lambda) = A_0 + A_{-1} \lambda^{-1}, \qquad \qquad \mathcal{B}(\lambda) = B_1 \lambda + B_0. \tag{8}$$

• Consider the matrices expressed in terms of the standard unit matrices  $E_{r,s} \in Mat_{n+1}(\mathbb{C})$  as

$$A_{0} = E_{1,n} + f_{0} E_{1,n+1} + E_{2,n+1}, \quad A_{-1} = \sum_{1 \le r \le n+1} \beta_{r} E_{r,r} + \sum_{1 \le r \le n} f_{r} E_{r+1,r} + \sum_{1 \le r \le n-1} E_{r+2,r},$$

$$B_1 = E_{1,n+1}, \quad B_0 = \sum_{1 \le r \le n+1} \frac{g_r}{g_r} E_{r,r} + \sum_{1 \le r \le n} E_{r+1,r}$$

• Let  $\alpha_0 = 1 + \beta_{n+1} - \beta_1$ ,  $\alpha_j = \beta_j - \beta_{j+1}$ ,  $j \in \mathbb{Z}/(n+1)\mathbb{Z} \setminus \{0\}$ .

**Theorem.** [Bobrova, 2024] There exists a set of the *g*-functions such that the compatibility condition of system (7) is equivalent to either the  $A_{2l}^{(1)}$  or  $A_{2l+1}^{(1)}$  system.

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### Bäcklund transformations and discrete dynamics

- Let the Cartan matrix C be of type  $A_n^{(1)}$ ,  $n \ge 2$  and  $I = \{0, 1, \dots, n\}$ .
- Let us set

$$\begin{aligned} s_i(\alpha_i) &= -\alpha_i, \qquad s_i(\alpha_j) = \alpha_j + \alpha_i \qquad (j = i \pm 1), \qquad s_i(\alpha_j) = \alpha_j \quad (j \neq i \pm 1), \\ s_i(f_i) &= f_i, \qquad s_i(f_j) = f_j \pm \alpha_i f_i^{-1} \qquad (j = i \pm 1), \qquad s_i(f_j) = f_j \qquad (j \neq i \pm 1), \\ \pi(\alpha_j) &= \alpha_{j+1}, \qquad \pi(f_j) = f_{j+1}, \qquad j \in \mathbb{Z}/(n+1)\mathbb{Z}. \end{aligned}$$

**Theorem.** [Bobrova, 2024] Transformations given above are Bäcklund transformations of the  $A_{2l}^{(1)}$  and  $A_{2l+1}^{(1)}$  systems. Moreover, they define a birational representation of the extended affine Weyl group of type  $A_n^{(1)}$ ,  $n \ge 2$ .

Note that the shift operators are given by

 $T_1 = \pi s_n s_{n-1} \dots s_1, \quad T_2 = s_1 \pi s_n \dots s_2, \quad \dots, \quad T_{n+1} = s_n \dots s_1 \pi.$  (9)

- They satisfy the relation  $T_1 T_2 \ldots T_{n+1} = 1$ .
- Thus, any n of them form a basis for the lattice and we can define a discrete system.

**Remark.** Cases n = 2 and n = 3 correspond to the P<sub>4</sub> and P<sub>5</sub> equations and discrete systems labeled by d-P( $E_6$ ) and d-P( $D_5$ ) respectively. For n = 1 one needs to consider the P<sub>2</sub> equation.

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d-Painlevé equations

### Overview

#### Agreements

- ▶ The elements q, p belong to  $\Re$  and all constant parameters labeled by greek letters  $\in \mathbb{C}$ .
- Usually, t is a central element, i.e.  $t \in \mathcal{I}(\mathcal{R})$ , except for the P<sub>2</sub> and P<sub>4</sub> systems.
- For discrete dynamics, we will use the standard notation. Namely, for a *T*-action of the translation operator *T*, we set  $T(f) = \overline{f}$  and  $T^{-1}(f) = \underline{f}$ .
- Regarding difference systems, we use  $T^n(f) = f_n$ .

#### Overview

- Thanks to the paper [Bershtein et al., 2023], we know that matrix Hamiltonian Painlevé systems of all types have Bäcklund transformations forming an affine Weyl group structure.
- We have reconstructed all the generators for the extended affine Weyl groups corresponding to the non-abelian Hamiltonian systems obtained in [Bobrova and Sokolov, 2023a].
- By using the translation operators as in [Sakai, 2001] or [Kajiwara et al., 2017], we have obtained the list of non-abelian discrete systems.
- ▶ They might be regarded as non-commutative analogs for the *d*-Painlevé systems.
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### P<sub>2</sub> case: the system and symmetries

► Consider the P<sub>2</sub> system [Retakh and Rubtsov, 2010] (see also [Adler and Sokolov, 2021])

$$\begin{cases} \dot{q} &= -q^2 + p - \frac{1}{2}t, \\ \dot{p} &= qp + pq + \alpha_1. \end{cases}$$
 P<sub>2</sub>

- Here we assume that t is also an element of  $\mathcal{R}$  such that  $\dot{t} = 1$ .
- Let  $\alpha_0 + \alpha_1 = 1$  and  $f := -p + 2q^2 + t$ .
- Its Bäcklund transformations are given below (cf. with [Bershtein et al., 2023])

	$\alpha_0$	$\alpha_1$	q	p	t
<i>s</i> 0	$-\alpha_0$	$\alpha_{1}+2\alpha_{0}$	$q - \alpha_0 f^{-1}$	$p - 2\alpha_0 q f^{-1} - 2\alpha_0 f^{-1} q + 2\alpha_0 f^{-2}$	t
<i>s</i> 1	$\alpha_0 + 2\alpha_1$	$-\alpha_1$	$q + \alpha_1 p^{-1}$	p	t
π	$\alpha_1$	$\alpha_0$	-q	$-p+2q^2+t$	t

These elements form an extended affine Weyl group of type A<sub>1</sub><sup>(1)</sup>:

$$\widetilde{W}(A_1^{(1)}) = \langle s_0, s_1; \pi \rangle,$$

$$s_i^2 = 1, \qquad \pi^2 = 1, \qquad \pi s_i = s_{i+1}\pi, \qquad i \in \mathbb{Z}/_{2\mathbb{Z}}.$$
(10)

### P<sub>2</sub> case: discrete dynamics

• Consider the translation operator  $T = s_1 \pi$ . It acts on the parameters according to the formula below and form a lattice on a line:

$$T(\alpha_0, \alpha_1) = (\alpha_0 - 1, \alpha_1 + 1).$$
(11)

The q and p variables change as follows

$$\bar{q} = s_1 \pi(q) = -s_1(q) = -q - \alpha_1 p^{-1}, \quad \bar{p} = s_1 \pi(p) = s_1(-p + 2q^2 + t) = -p + 2\bar{q}^2 + t.$$

So, we obtain the system

$$\bar{\alpha}_0 = \alpha_0 - 1, \qquad \bar{\alpha}_1 = \alpha_1 + 1,$$
  
$$\bar{q} + q = -\alpha_1 p^{-1}, \qquad \bar{p} + p = t + 2\bar{q}^2.$$
  
$$\mathsf{d}\text{-P}(E_7)$$

- ▶ It generalizes to the non-commutative case the d-P( $E_7$ ) equation from [Sakai, 2001] (p. 206).
- The d-P( $E_7$ ) system can be rewritten in the difference form

$$\begin{cases} q_{n+1} + q_n = -\alpha_{1,n} p_n^{-1} \\ p_n + p_{n-1} = 2q_n^2 + t, \end{cases} \qquad \alpha_{1,n} = \alpha_1 + n, \tag{12}$$

which reduces to the following second-order difference equation:

$$\alpha_{1,n} (q_{n+1}+q_n)^{-1} + \alpha_{1,n-1} (q_n+q_{n-1})^{-1} = -2q_n^2 - t, \quad \alpha_{1,n} = \alpha_1 + n.$$
 alt-d-P<sub>1</sub>

 $d-P(E_7)$  system (1)

Lax pair

► One may also consider the corresponding discrete linear problem

$$\begin{cases} \partial_{\lambda} Y_n = \mathcal{A}_n Y_n, \\ Y_{n+1} = \mathcal{B}_n Y_n. \end{cases}$$
(13)

• A Lax pair for the  $d-P(E_7)$  is given by

$$\begin{split} \mathcal{A}_n &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \lambda^2 + \begin{pmatrix} 0 & 1 \\ 2p & 0 \end{pmatrix} \lambda + \begin{pmatrix} -p + \frac{1}{2}t & -q \\ 2pq + 2\alpha_1 & p - \frac{1}{2}t \end{pmatrix}, \\ \mathcal{B}_n &= \begin{pmatrix} -2 & 0 \\ 0 & 0 \end{pmatrix} \lambda + \begin{pmatrix} -2q & -1 \\ -2\bar{p} & 0 \end{pmatrix}, \end{split}$$

where  $t, \lambda \in \mathcal{Z}(\mathcal{R})$ .

Note that the compatibility condition is satisfied, since the commutator [p, q] is invariant under the map

$$\psi: \mathbb{R}^2 \to \mathbb{R}^2, \quad (q, p) \mapsto (\bar{q}, \bar{p}) = \left(-p + t + 2q^2, -q - \bar{\alpha}_1(-p + t + 2q^2)^{-1}\right). \tag{14}$$

• Once  $t \in \mathbb{R}$ , the commutator [p, q] is no longer a conserved quantity.

**Remark.** The latter fact might have been caused by the Hamiltonian structure similarly to the case of non-abelian Hamiltonian ODEs (see Lemma 1 in [Bobrova and Sokolov, 2023a] and its generalisation, Lemma 2.1, in [Bobrova, 2023]).

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# $d-P(E_7)$ system (2)

#### Hamiltonian structure

- Let us use non-abelian partial derivatives introduced in [Kontsevich, 1993].
- Similar to [Mase et al., 2020], we call a difference discrete system Hamiltonian if there exists a function  $H = H(q, \bar{p})$  such that the system can be rewritten in the form

$$p = \partial_q H, \qquad \bar{q} = \partial_{\bar{p}} H.$$
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For the d-P( $E_7$ ) system, a Hamiltonian is  $H = -q \,\bar{p} + tq + \frac{1}{3}q^3 - \bar{\alpha}_1 \ln \bar{p}$ , where for the symbol ln f we define the right logarithmic derivative by  $d_t(\ln f) := f^{-1} \dot{f}$ .

Then, the non-abelian derivatives are

$$\partial_q H = -\bar{p} + q^2 + t, \qquad \partial_{\bar{p}} H = -q - \bar{\alpha}_1 \, \bar{p}^{-1}, \qquad \partial_t H = q$$
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and (15) is equivalent to the  $d-P(E_7)$  system.

#### **Continuous limit**

One may consider a non-abelian analog for the continuous limit. Then, by using the formulas

$$q = 1 + \varepsilon^2 Q - \frac{1}{6} \varepsilon^3 P, \quad p = -2 + 2\varepsilon^2 Q + \frac{2}{3} \varepsilon^3 P, \quad t = -6 + \frac{1}{3} \varepsilon^4 T, \quad \alpha_1 = 4 + \frac{2}{3} \varepsilon^4 T,$$

the d-P( $E_7$ ) has the P<sub>1</sub> system in the limit  $\varepsilon \rightarrow 0$ :

$$\dot{q} = p,$$
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### **Further questions**

#### A study of the obtained *d*-Painlevé equations

- ▶ We expect that these equations admit Lax pairs and have a Hamiltonian form.
- We also expect that they have a continuous limit to known non-abelian differential Painlevé equations obtained in [Bobrova and Sokolov, 2023b].
- Commutative d- and differential Painlevé equations are connected with the orthogonal polynomials [Van Assche, 2022]. Orthogonal polynomials have a non-commutative analog (see [Gelfand et al., 1995]). We assume that our equations are connected with them.

#### Other discrete Painlevé equations

- Our method can be applied to the *q*-discrete Painlevé equations.
- In particular, one may define a non-abelian version for the q-P<sub>6</sub> equation which generalizes the matrix equation obtained in [Kawakami, 2020] to the purely non-abelian case. (an ongoing project)
- ▶ How to derive a non-commutative *ell*-discrete Painlevé equation?

#### Non-abelian geometry related to Painlevé equations

- ▶ What is the Okamoto space of initial data of non-abelian differential Painlevé equations?
- We would like to generalize the method of the Painlevé equations' classification introduced in the Sakai's paper [Sakai, 2001]. Recent developments might be found in [Rains, 2019].

#### Cluster algebras and discrete Painlevé equations

It is known that discrete Painlevé equations are connected with claster algebras (see, e.g. [Bershtein et al., 2018]). Might we have the same connection in the non-abelian case?

# Many thanks!

### **References** I

[Ablowitz et al., 1980] Ablowitz, M. J., Ramani, A., and Segur, H. (1980). A connection between nonlinear evolution equations and ordinary differential equations of P-type. II. Journal of Mathematical Physics, 21(5):1006–1015.

[Adler, 1994] Adler, V. E. (1994). Nonlinear chains and Painlevé equations. Physica D: Nonlinear Phenomena, 73(4):335–351.

[Adler, 2020] Adler, V. E. (2020).

Painlevé type reductions for the non-Abelian Volterra lattices. Journal of Physics A: Mathematical and Theoretical, 54(3):035204. arXiv:2010.09021.

[Adler and Sokolov, 2021] Adler, V. E. and Sokolov, V. V. (2021).
 On matrix Painlevé II equations.
 Theoret. and Math. Phys., 207(2):188–201.
 arXiv:2012.05639.

[Bershtein et al., 2018] Bershtein, M., Gavrylenko, P., and Marshakov, A. (2018). Cluster integrable systems, q-Painlevé equations and their quantization. Journal of High Energy Physics, 2018(2):1 – 29. arXiv:1711.02063.

[Bershtein et al., 2023] Bershtein, M., Grigorev, A., and Shchechkin, A. (2023). Hamiltonian reductions in matrix Painlevé systems. Letters in Mathematical Physics, 113(2):47. arXiv:2208.04824.

### **References II**

[Bertola et al., 2018] Bertola, M., Cafasso, M., and Rubtsov, V. (2018). Noncommutative Painlevé equations and systems of Calogero type. Communications in Mathematical Physics, 363(2):503–530. arXiv:1710.00736.

[Bobrova, 2024] Bobrova, I. (2024).

Affine Weyl groups and non-Abelian discrete systems: an application to the *d*-Painlevé equations. arXiv preprint arXiv:2403.18463.

[Bobrova, 2023] Bobrova, I. (2023). Équations de Painlevé non abéliennes. PhD thesis, Reims.

[Bobrova et al., 2023] Bobrova, I., Retakh, V., Rubtsov, V., and Sharygin, G. (2023). Non-Abelian discrete Toda chains and related lattices. *Physica D: Nonlinear Phenomena (under review)*. arXiv:2311.11124.

[Bobrova and Sokolov, 2023a] Bobrova, I. and Sokolov, V. (2023a). Classification of Hamiltonian non-abelian Painlevé type systems. *Journal of Nonlinear Mathematical Physics*, 30:646–662. arXiv:2209.00258.

[Bobrova and Sokolov, 2023b] Bobrova, I. and Sokolov, V. (2023b). On classification of non-abelian Painlevé type systems. Journal of Geometry and Physics, 191:104885. arXiv:2303.10347.

### **References III**

[Cafasso and Manuel, 2014] Cafasso, M. and Manuel, D. (2014). Non-commutative Painlevé equations and Hermite-type matrix orthogonal polynomials. *Communications in Mathematical Physics*, 326(2):559–583. arXiv:1301.2116.

[Cafasso et al., 2018] Cafasso, M., Manuel, D., et al. (2018).

The Toda and Painlevé systems associated with semiclassical matrix-valued orthogonal polynomials of Laguerre type. SIGMA. Symmetry, Integrability and Geometry: Methods and Applications, 14:076.

arXiv:1801.08740.

[Cassatella-Contra et al., 2014] Cassatella-Contra, G. A., Manas, M., and Tempesta, P. (2014). Singularity confinement for matrix discrete Painlevé equations. Nonlinearity, 27(9):2321.

[Dzhamay and Takenawa, 2018] Dzhamay, A. and Takenawa, T. (2018). On some applications of Sakai's geometric theory of discrete Painlevé equations. SIGMA. Symmetry, Integrability and Geometry: Methods and Applications, 14:075. arXiv:1804.10341.

[Gambier, 1910] Gambier, B. (1910).

Sur les équations différentielles du second ordre et du premier degré dont l'intégrale générale est à points critiques fixes.

Acta Mathematica, 33(1):1-55.

### **References IV**

[Gelfand et al., 1995] Gelfand, I. M., Krob, D., Lascoux, A., Leclerc, B., Retakh, V. S., and Thibon, J.-Y. (1995). Noncommutative Symmetric Functions. Advances in Mathematics, 2(112):218–348. arXiv:hep-th/9407124.

[Joshi, 2020] Joshi, N. (2020). Discrete Painlevé equations. Notices of the American Mathematical Society, 67(6):797–805.

[Kajiwara et al., 2017] Kajiwara, K., Noumi, M., and Yamada, Y. (2017). Geometric aspects of Painlevé equations. Journal of Physics A: Mathematical and Theoretical, 50(7):073001. arXiv:1509.08186.

[Kawakami, 2015] Kawakami, H. (2015). Matrix Painlevé systems. Journal of Mathematical Physics, 56(3):033503.

[Kawakami, 2020] Kawakami, H. (2020).

A q-analogue of the matrix sixth Painlevé system. Journal of Physics A: Mathematical and Theoretical, 53(49):495203. arXiv:2301.12837.

[Kontsevich, 1993] Kontsevich, M. (1993).

Formal (non)-commutative symplectic geometry, The Gelfand Mathematical Seminars, 1990–1992. Fields Institute Communications, Birkhäuser Boston, pages 173–187.

### **References V**

[Mase et al., 2020] Mase, T., Nakamura, A., and Sakai, H. (2020).
 Discrete Hamiltonians of discrete Painlevé equations.
 In Annales de la Faculté des sciences de Toulouse: Mathématiques, volume 29 (5), pages 1251–1264.

[Nagoya, 2004] Nagoya, H. (2004).

Quantum Painlevé systems of type  $A_l^{(1)}$ . International Journal of Mathematics, 15(10):1007–1031. arXiv:math/0402281v2.

[Noumi and Yamada, 1998] Noumi, M. and Yamada, Y. (1998). Affine Weyl groups, discrete dynamical systems and Painlevé equations. Communications in Mathematical Physics, 199:281–295.

[Noumi and Yamada, 2000] Noumi, M. and Yamada, Y. (2000). Affine Weyl group symmetries in Painlevé type equations. *Citeseer.* 

[Olver and Sokolov, 1998] Olver, P. J. and Sokolov, V. V. (1998). Integrable evolution equations on associative algebras. Communications in Mathematical Physics, 193(2):245–268.

[Olver and Wang, 2000] Olver, P. J. and Wang, J. P. (2000). Classification of integrable one-component systems on associative algebras. *Proceedings of the London Mathematical Society*, 81(3):566–586.

### **References VI**

[Painlevé, 1902] Painlevé, P. (1902).

Sur les équations différentielles du second ordre et d'ordre supérieur dont l'intégrale générale est uniforme. Acta mathematica, 25:1–85.

[Rains, 2019] Rains, E. M. (2019). The birational geometry of noncommutative surfaces. arXiv preprint arXiv:1907.11301.

[Retakh and Rubtsov, 2010] Retakh, V. S. and Rubtsov, V. N. (2010). Noncommutative Toda Chains, Hankel Quasideterminants and Painlevé II Equation. Journal of Physics A, Mathematical and Theoretical, 43(50):505204. arXiv:1007.4168.

[Sakai, 2001] Sakai, H. (2001).

Rational surfaces associated with affine root systems and geometry of the Painlevé equations. Communications in Mathematical Physics, 220(1):165–229.

[Van Assche, 2022] Van Assche, W. (2022). Orthogonal polynomials, Toda lattices and Painlevé equations. *Physica D: Nonlinear Phenomena*, 434:133214. arXiv:2202.11017.

[Veselov and Shabat, 1993] Veselov, A. P. and Shabat, A. B. (1993). Dressing chains and the spectral theory of the Schrödinger operator. Functional Analysis and Its Applications, 27(2):81–96.

### **References VII**

[Wadati and Kamijo, 1974] Wadati, M. and Kamijo, T. (1974).

On the extension of inverse scattering method.

Progress of theoretical Physics, 52(2):397-414.

[Zakharov and Shabat, 1974] Zakharov, V. E. and Shabat, A. B. (1974).

A scheme for integrating the nonlinear equations of mathematical physics by the method of the inverse scattering problem. I.

Functional Analysis and Its Applications, 8(3):43-53.